

## $\mathbb{R}$ -Complex Finsler Spaces with an Arctangent Finsler Metric

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**Abstract.** In this paper, we have defined the concept of the  $\mathbb{R}$ -complex Finsler space with an arctangent  $(\alpha, \beta)$ -metric  $F = \alpha + \epsilon\beta + \beta \tan^{-1}(\beta/\alpha)$ . For this metric, we have obtained the fundamental metric tensor fields  $g_{ij}$  and  $g_{i\bar{j}}$  as well as their determinants and inverse tensor fields. Further, some properties of non-Hermitian  $\mathbb{R}$ -complex Finsler spaces with this metric have been described.

**Keywords:** Complex Finsler space,  $\mathbb{R}$ -complex Finsler space, Fundamental metric tensors.

### 1. Introduction

The concept of  $(\alpha, \beta)$ -metric has multiple applications in Physics, Ecology, and Biology [7, 15]. It was investigated first by M. Matsumoto in perspective of the generalization of Rander's metric [9]. Afterward, many geometers studied it in great detail and have derived different kinds of  $(\alpha, \beta)$ -metrics like the general  $(\alpha, \beta)$ -metric, Kropina metric, Einstein metric, Matsumoto metric, and Exponential metric etc. in some different geometrical points of view.

The theories of  $\mathbb{R}$ -complex Finsler spaces are very new and it was introduced first by G. D. Rizza [13]. G. Munteanu and M. Purcaru [11] have extended the idea of the complex Finsler spaces [1, 3, 10] and got another class of such

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spaces. N. Aldea [4] had investigated a class of complex Finsler spaces with two-dimensions. Recently, many researchers [2, 5, 8] have obtained many fundamental results on  $\mathbb{R}$ -Complex Finsler spaces.

The paper follows ideas from real Finsler space with an arctangent metric and introduces the similar notion on  $\mathbb{R}$ -Complex Finsler Spaces with it defined by

$$F = \alpha + \epsilon\beta + \beta \tan^{-1}\left(\frac{\beta}{\alpha}\right). \quad (1.1)$$

For the above mentioned metric, we have obtained the fundamental metric tensor fields  $g_{ij}$  and  $g_{i\bar{j}}$  along with their determinants and the inverse tensor fields. Further, we have discussed some properties of non-Hermitian  $\mathbb{R}$ -complex Finsler spaces with the metric given in equation (1.1).

## 2. $\mathbb{R}$ -Complex Finsler Spaces

Let  $M$  be a complex manifold with  $\dim_c M = n$ ,  $(z^k)$  be local complex coordinates in a chart  $(U, \phi)$  and  $T'M$  its holomorphic tangent bundle. It has a natural structure of complex manifold,  $\dim_c T'M = 2n$  and the induced coordinates in a local chart on  $u \in T'M$  are denoted by  $u = (z^k, \eta^k)$ . The changes of local coordinates in  $u$  are given by the rules:

$$z'^k = z'^k(z), \quad \eta'^k = \frac{\partial z'^k}{\partial z^j} \eta^j. \quad (2.1)$$

The natural frame  $\left\{ \frac{\partial}{\partial z^k}, \frac{\partial}{\partial \eta^k} \right\}$  of  $T'_u(T'M)$  transforms with the Jacobi matrix of equation (2.1) changes,

$$\frac{\partial}{\partial z^k} = \frac{\partial z'^j}{\partial z^k} \frac{\partial}{\partial z'^j} + \frac{\partial^2 z'^j}{\partial z^k \partial z^h} \eta^h \frac{\partial}{\partial \eta'^j}, \quad \frac{\partial}{\partial \eta^k} = \frac{\partial z'^j}{\partial z^k} \frac{\partial}{\partial \eta'^j}.$$

A complex non-linear connection, briefly c.n.c, is a supplementary distribution  $H(T'M)$  to the vertical distribution  $V(T'M)$  in  $T'(T'M)$ . The vertical distribution is spanned by  $\left\{ \frac{\partial}{\partial \eta^k} \right\}$  and an adapted frame in  $H(T'M)$  is

$$\frac{\delta}{\delta z^k} = \frac{\partial}{\partial z^k} - N_k^j \frac{\partial}{\partial \eta^j},$$

where  $N_k^j$  are the coefficients of the c.n.c. and they have a certain rule of change at (2.1) so that  $\frac{\delta}{\delta z^k}$  transform like vectors on the base manifold  $M$ . Next, we use the abbreviations:

$$\partial_k = \frac{\partial}{\partial z^k}, \quad \delta_k = \frac{\delta}{\delta z^k}, \quad \dot{\partial}_k = \frac{\partial}{\partial \eta^k}, \quad \partial_{\bar{k}}, \delta_{\bar{k}}, \dot{\partial}_{\bar{k}}$$

for their conjugates. The dual adapted basis of  $\{\delta_k, \dot{\partial}_k\}$  are  $\{dz^k, \delta\eta^k = d\eta^k + N_j^k dz^j\}$  and  $\{d\bar{z}^k, \delta\bar{\eta}^k\}$  their conjugates.

**Definition 2.1.** [11] An  $\mathbb{R}$ -Complex Finsler metric on  $M$  is a continuous function  $F : T'M \rightarrow \mathbb{R}_+$  satisfying:

- i)  $L := F^2$  is smooth on  $\widetilde{T'M}$  (except the 0 sections),
- ii)  $F(z, \eta) \geq 0$ , the equality holds if and only if  $\eta = 0$ ,
- iii)  $F(z, \lambda\eta, \bar{z}, \lambda\bar{\eta}) = |\lambda| F(z, \eta, \bar{z}, \bar{\eta}), \forall \lambda \in \mathbb{R}$ .

Using assertion (i) and (iii) of the definition 2.1,  $L$  is  $(2, 0)$  homogeneous with respect to the real scalars  $\lambda$ , i.e. ,  $L(z, \lambda\eta, \bar{z}, \lambda\bar{\eta}) = \lambda^2 L(z, \eta, \bar{z}, \bar{\eta}), \lambda \in \mathbb{R}$ .

**Definition 2.2.** [6] An  $\mathbb{R}$ -Complex Finsler spaces with  $(\alpha, \beta)$ -metric is a pair  $(M, F)$ , where the fundamental function  $F(z, \eta, \bar{z}, \bar{\eta})$  is  $\mathbb{R}$ -homogeneous by means of functions  $\alpha(z, \eta, \bar{z}, \bar{\eta})$  and  $\beta(z, \eta, \bar{z}, \bar{\eta})$ ,

$$F(z, \eta, \bar{z}, \bar{\eta}) = F(\alpha(z, \eta, \bar{z}, \bar{\eta}), \beta(z, \eta, \bar{z}, \bar{\eta})), \quad (2.2)$$

where

$$\begin{aligned} \alpha^2(z, \eta, \bar{z}, \bar{\eta}) &= \frac{1}{2}(a_{ij}\eta^i\eta^j + 2a_{i\bar{j}}\eta^i\bar{\eta}^j + a_{\bar{i}j}\bar{\eta}^i\eta^j) \\ &= \text{Re}\{a_{ij}\eta^i\eta^j + a_{i\bar{j}}\eta^i\bar{\eta}^j\}, \\ \beta(z, \eta, \bar{z}, \bar{\eta}) &= \frac{1}{2}(b_i\eta^i + b_{\bar{i}}\bar{\eta}^i) \\ &= \text{Re}\{b_i\eta^i\}, \end{aligned}$$

with  $a_{ij} = a_{ij}(z)$ ,  $a_{i\bar{j}} = a_{i\bar{j}}(z)$  and  $b_i = b_i(z)$ ,  $b_i(z)dz^i$  is a  $(1, 0)$ -differential form on complex manifold  $M$ .

If  $a_{ij} = 0$  and  $a_{i\bar{j}}$  invertible, then the space is said to be of *Hermitian space*. If  $a_{i\bar{j}} = 0$  and  $a_{ij}$  invertible, then the space is called *non-Hermitian space*.

Indeed,  $\alpha$  and  $\beta$  are homogeneous with respect to  $\eta$  and  $\bar{\eta}$ , i.e.  $\alpha(z, \lambda\eta, \bar{z}, \lambda\bar{\eta}) = \lambda\alpha(z, \eta, \bar{z}, \bar{\eta})$  and  $\beta(z, \lambda\eta, \bar{z}, \lambda\bar{\eta}) = \lambda\beta(z, \eta, \bar{z}, \bar{\eta})$ , for any  $\lambda \in \mathbb{R}_+$ . Since  $L$  is  $(2, 0)$  homogeneous with respect to  $\lambda$ , by using the homogeneity property following equalities hold [6]:

$$\begin{aligned} \alpha L_\alpha + \beta L_\beta &= 2L, \\ \alpha L_{\alpha\alpha} + \beta L_{\alpha\beta} &= L_\alpha, \\ \alpha L_{\alpha\beta} + \beta L_{\beta\beta} &= L_\beta, \\ \alpha^2 L_{\alpha\alpha} + 2\alpha\beta L_{\alpha\beta} + \beta^2 L_{\beta\beta} &= 2L, \\ \frac{\partial \alpha}{\partial \eta^i} \eta^i + \frac{\partial \alpha}{\partial \bar{\eta}^j} \bar{\eta}^j &= \alpha, \\ \frac{\partial \beta}{\partial \eta^i} \eta^i + \frac{\partial \beta}{\partial \bar{\eta}^j} \bar{\eta}^j &= \beta, \end{aligned} \quad (2.3)$$

where

$$L_\alpha = \frac{\partial L}{\partial \alpha}, \quad L_\beta = \frac{\partial L}{\partial \beta}, \quad L_{\alpha\alpha} = \frac{\partial^2 L}{\partial \alpha^2}, \quad L_{\beta\beta} = \frac{\partial^2 L}{\partial \beta^2}, \quad L_{\alpha\beta} = \frac{\partial^2 L}{\partial \alpha \partial \beta}.$$

We consider

$$\frac{\partial \alpha}{\partial \eta^i} = \frac{1}{2\alpha}(a_{ij}\eta^j + a_{i\bar{j}}\bar{\eta}^j) = \frac{1}{2\alpha}l_i, \quad \frac{\partial \beta}{\partial \eta^i} = \frac{1}{2}b_i,$$

and

$$\begin{aligned} \eta_i &= \frac{\partial L}{\partial \eta^i} = \frac{\partial}{\partial \eta^i} F^2 = 2F \frac{\partial F}{\partial \eta^i} \\ &= \rho_0 l_i + \rho_1 b_i, \end{aligned}$$

where

$$l_i = a_{ij}\eta^j + a_{i\bar{j}}\bar{\eta}^j, \quad (2.4)$$

$$b_l = a_{kl}b^k + a_{l\bar{k}}\bar{b}^k, \quad (2.5)$$

$$\rho_0 = \frac{1}{2} \frac{L_{\alpha\alpha}}{\alpha}, \quad \rho_1 = \frac{1}{2} L_{\beta\beta}. \quad (2.6)$$

Differentiating  $\rho_0$  and  $\rho_1$  w.r.t.  $\eta^j$ , we get

$$\begin{aligned} \frac{\partial \rho_0}{\partial \eta^j} &= \rho_{-2} l_j + \rho_{-1} b_j, \\ \frac{\partial \rho_1}{\partial \eta^i} &= \rho_{-1} l_i + \mu_0 b_i, \end{aligned}$$

where

$$\rho_{-2} = \frac{\alpha L_{\alpha\alpha} - L_{\alpha}}{4\alpha^3}, \quad \rho_{-1} = \frac{L_{\alpha\beta}}{4\alpha}, \quad \mu_0 = \frac{L_{\beta\beta}}{4}. \quad (2.7)$$

The quantities  $\rho_{-2}, \rho_{-1}, \rho_0, \rho_1, \mu_0$  are the invariants of the  $\mathbb{R}$ -complex Finsler space with  $(\alpha, \beta)$ -metric [14].

In [11], an  $\mathbb{R}$ -Complex Finsler space the following conditions hold:

$$\begin{aligned} \frac{\partial L}{\partial \eta^i} \eta^i + \frac{\partial L}{\partial \bar{\eta}^i} \bar{\eta}^i &= 2L; \quad g_{ij} \eta^i + g_{j\bar{i}} \bar{\eta}^i = \frac{\partial L}{\partial \eta^j}; \\ \frac{\partial g_{ik}}{\partial \eta^j} \eta^j + \frac{\partial g_{i\bar{k}}}{\partial \bar{\eta}^j} \bar{\eta}^j &= 0; \quad \frac{\partial g_{i\bar{k}}}{\partial \eta^j} \eta^j + \frac{\partial g_{i\bar{k}}}{\partial \bar{\eta}^j} \bar{\eta}^j = 0; \\ 2L &= g_{ij} \eta^i \eta^j + g_{i\bar{j}} \bar{\eta}^i \bar{\eta}^j + 2g_{i\bar{j}} \eta^i \bar{\eta}^j, \end{aligned}$$

where

$$g_{ij} = \frac{\partial^2 L}{\partial \eta^i \partial \eta^j}, \quad g_{i\bar{j}} = \frac{\partial^2 L}{\partial \eta^i \partial \bar{\eta}^j}, \quad \text{and} \quad g_{\bar{i}j} = \frac{\partial^2 L}{\partial \bar{\eta}^i \partial \eta^j},$$

are the metric tensors of space.

**Theorem 2.3.** [6] *The metric tensor fields of  $\mathbb{R}$ -complex Finsler space with  $(\alpha, \beta)$ -metric are given by*

$$\begin{aligned} g_{ij} &= \rho_0 a_{ij} + \rho_{-2} l_i l_j + \mu_0 b_i b_j + \rho_{-1} (b_j l_i + b_i l_j), \\ g_{i\bar{j}} &= \rho_0 a_{i\bar{j}} + \rho_{-2} l_i l_{\bar{j}} + \mu_0 b_i b_{\bar{j}} + \rho_{-1} (b_{\bar{j}} l_i + b_i l_{\bar{j}}), \end{aligned} \quad (2.8)$$

where the quantities  $\rho_{-2}, \rho_{-1}, \rho_0, \rho_1, \mu_0$  are defined in the symbols of equations (2.6) and (2.7).

For obtaining the inverse and determinant of the tensor field  $g_{ij}$ , one can follow the following proposition:

**Proposition 2.4.** [5] *Suppose*

- $(Q_{ij})$  is a non-singular  $n \times n$  complex matrix with inverse  $(Q^{ji})$ ;
- $C_i$  and  $C_{\bar{i}} := \bar{C}_i, i = 1, \dots, n$  are complex numbers;
- $C^i := Q^{ji}C_j$  and  $C_i$  are conjugates to each other;  $C^2 := C^iC_i = \bar{C}^iC_{\bar{i}}; H_{ij} := Q_{ij} \pm C_iC_j$ .

Then

- i)  $\det(H_{ij}) = (1 \pm C^2)\det(Q_{ij})$ ,
- ii) whenever  $1 \pm C^2 \neq 0$ , the matrix  $(H_{ij})$  is invertible and its inverse is

$$H^{ji} = Q^{ji} \mp \frac{1}{1 \pm C^2} C^i C^j.$$

### 3. $\mathbb{R}$ -Complex Finsler Space with an Arctangent Metric

An  $\mathbb{R}$ -Complex Finsler spaces  $(M, F)$  is known as  $\mathbb{R}$ -Complex arctangent Finsler space if  $F$  satisfies the equation (2.2).

From the definition 2.1 (i), we have

$$L(\alpha, \beta) = \{\alpha + \epsilon\beta + \beta \tan^{-1}(\beta/\alpha)\}^2. \quad (3.1)$$

From above equation, we get

$$\begin{aligned} L_\alpha &= \frac{2\alpha^2}{\alpha^2 + \beta^2} \{\alpha + \epsilon\beta + \beta \tan^{-1}(\beta/\alpha)\}, \\ L_{\alpha\alpha} &= \frac{2\alpha}{(\alpha^2 + \beta^2)^2} [2\beta^3\{\epsilon + \tan^{-1}(\beta/\alpha)\} + \alpha(\alpha^2 + 2\beta^2)], \\ L_{\alpha\beta} &= \frac{2\alpha^2}{(\alpha^2 + \beta^2)^2} [(\alpha^2 - \beta^2)\{\epsilon + \tan^{-1}(\beta/\alpha)\} - \beta\alpha], \\ L_\beta &= \frac{2}{\alpha^2 + \beta^2} \{\alpha + \epsilon\beta + \beta \tan^{-1}(\beta/\alpha)\} [(\alpha^2 + \beta^2)\{\epsilon + \tan^{-1}(\beta/\alpha)\} + \beta\alpha], \\ L_{\beta\beta} &= \frac{2}{(\alpha^2 + \beta^2)^2} \left[ [(\alpha^2 + \beta^2)\{\epsilon + \tan^{-1}(\beta/\alpha)\} + \beta\alpha]^2 \right. \\ &\quad \left. + 2\alpha^3\{\alpha + \epsilon\beta + \beta \tan^{-1}(\beta/\alpha)\} \right]. \end{aligned} \quad (3.2)$$

Substituting  $L_\alpha$ ,  $L_{\alpha\alpha}$ ,  $L_\beta$ ,  $L_{\beta\beta}$ , and  $L_{\alpha\beta}$  from above in the system of equations (2.3), we get

$$\alpha L_\alpha + \beta L_\beta = 2 \{ \alpha + \epsilon \beta + \beta \tan^{-1}(\beta/\alpha) \}^2 = 2L, \quad (3.3)$$

$$\alpha L_{\alpha\alpha} + \beta L_{\alpha\beta} = \frac{2\alpha^2}{\alpha^2 + \beta^2} \{ \alpha + \epsilon \beta + \beta \tan^{-1}(\beta/\alpha) \} = L_\alpha, \quad (3.4)$$

$$\alpha L_{\alpha\beta} + \beta L_{\beta\beta} = L_\beta,$$

$$\alpha^2 L_{\alpha\alpha} + 2\alpha\beta L_{\alpha\beta} + \beta^2 L_{\beta\beta} = 2L.$$

In the same way, one can verify the rest equalities of the system of equations (2.3).

**Proposition 3.1.** *The invariants of an  $\mathbb{R}$ -Complex Finsler space  $(M, F)$ , where  $F$  is an arctangent metric, are given in the system of equations:*

*Now, using the equations (2.6), (2.7), and (3.2), we get*

$$\begin{aligned} \rho_0 &= \frac{\alpha}{\alpha^2 + \beta^2} \{ \alpha + \epsilon \beta + \beta \tan^{-1}(\beta/\alpha) \}, \\ \rho_1 &= \frac{1}{\alpha^2 + \beta^2} \{ \alpha + \epsilon \beta + \beta \tan^{-1}(\beta/\alpha) \} \left[ (\alpha^2 + \beta^2) \{ \epsilon + \tan^{-1}(\beta/\alpha) \} + \beta \alpha \right], \\ \rho_{-2} &= \frac{-\beta}{2\alpha (\alpha^2 + \beta^2)^2} \left[ (\alpha^2 - \beta^2) \{ \epsilon + \tan^{-1}(\beta/\alpha) \} - \beta \alpha \right], \\ \rho_{-1} &= \frac{\alpha}{2(\alpha^2 + \beta^2)^2} \left[ (\alpha^2 - \beta^2) \{ \epsilon + \tan^{-1}(\beta/\alpha) \} - \beta \alpha \right], \\ \mu_0 &= \frac{1}{2(\alpha^2 + \beta^2)^2} \left[ [(\alpha^2 + \beta^2) \{ \epsilon + \tan^{-1}(\beta/\alpha) \} + \beta \alpha]^2 \right. \\ &\quad \left. + 2\alpha^3 \{ \alpha + \epsilon \beta + \beta \tan^{-1}(\beta/\alpha) \} \right]. \end{aligned} \quad (3.5)$$

**Theorem 3.2.** *The metric tensor fields of an  $\mathbb{R}$ -Complex Finsler space  $(M, F)$ , where  $F$  is an arctangent metric, are given in equations:*

*Now, using the invariants given in equation (3.5) and theorem 2.3, we get*

$$\begin{aligned} g_{ij} &= \frac{\alpha}{\alpha^2 + \beta^2} \{ \alpha + \epsilon \beta + \beta \tan^{-1}(\beta/\alpha) \} a_{ij} \\ &\quad + \frac{-\beta}{2\alpha (\alpha^2 + \beta^2)^2} \left[ (\alpha^2 - \beta^2) \{ \epsilon + \tan^{-1}(\beta/\alpha) \} - \beta \alpha \right] l_i l_j \\ &\quad + \frac{1}{2(\alpha^2 + \beta^2)^2} \left[ [(\alpha^2 + \beta^2) \{ \epsilon + \tan^{-1}(\beta/\alpha) \} + \beta \alpha]^2 \right. \\ &\quad \left. + 2\alpha^3 \{ \alpha + \epsilon \beta + \beta \tan^{-1}(\beta/\alpha) \} \right] b_i b_j \\ &\quad + \frac{\alpha}{2(\alpha^2 + \beta^2)^2} \left[ (\alpha^2 - \beta^2) \{ \epsilon + \tan^{-1}(\beta/\alpha) \} - \beta \alpha \right] (b_j l_i + b_i l_j), \end{aligned} \quad (3.6)$$

and

$$\begin{aligned}
 g_{i\bar{j}} &= \frac{\alpha}{\alpha^2 + \beta^2} \left\{ \alpha + \epsilon\beta + \beta \tan^{-1}(\beta/\alpha) \right\} a_{i\bar{j}} \\
 &\quad - \frac{\beta}{2\alpha(\alpha^2 + \beta^2)^2} \left[ (\alpha^2 - \beta^2) \{ \epsilon + \tan^{-1}(\beta/\alpha) \} - \beta\alpha \right] l_i l_{\bar{j}} \\
 &\quad + \frac{1}{2(\alpha^2 + \beta^2)^2} \left[ (\alpha^2 + \beta^2) \{ \epsilon + \tan^{-1}(\beta/\alpha) \} + \beta\alpha \right]^2 \\
 &\quad + 2\alpha^3 \{ \alpha + \epsilon\beta + \beta \tan^{-1}(\beta/\alpha) \} b_i b_{\bar{j}} \\
 &\quad + \frac{\alpha}{2(\alpha^2 + \beta^2)^2} \left[ (\alpha^2 - \beta^2) \{ \epsilon + \tan^{-1}(\beta/\alpha) \} - \beta\alpha \right] (b_{\bar{j}} l_i + b_i l_{\bar{j}}). \quad (3.7)
 \end{aligned}$$

Or

the equations (3.6) and (3.7) can be written in the following equivalent forms:

$$g_{ij} = \rho_0(a_{ij} - t_1 l_i l_j + t_2 b_i b_j + t_3 \eta_i \eta_j), \quad (3.8)$$

$$g_{i\bar{j}} = \rho_0(a_{i\bar{j}} - t_1 l_i l_{\bar{j}} + t_2 b_i b_{\bar{j}} + t_3 \eta_i \eta_{\bar{j}}), \quad (3.9)$$

where

$$\begin{aligned}
 t_1 &= \frac{[(\alpha^2 - \beta^2) \{ \epsilon + \tan^{-1}(\beta/\alpha) \} - \beta\alpha]}{2\alpha^2(\alpha^2 + \beta^2)[(\alpha^2 + \beta^2) \{ \epsilon + \tan^{-1}(\beta/\alpha) \} + \beta\alpha]}, \\
 t_2 &= \frac{\{ \alpha + \epsilon\beta + \beta \tan^{-1}(\beta/\alpha) \}}{\alpha}, \\
 t_3 &= \frac{(\alpha^2 + \beta^2)[(\alpha^2 - \beta^2) \{ \epsilon + \tan^{-1}(\beta/\alpha) \} - \beta\alpha]}{2\alpha \{ \alpha + \epsilon\beta + \beta \tan^{-1}(\beta/\alpha) \}^3 [(\alpha^2 + \beta^2) \{ \epsilon + \tan^{-1}(\beta/\alpha) \} + \beta\alpha]}. \quad (3.10)
 \end{aligned}$$

*Proof.* Using the relations (3.5) in theorem 2.3 by direct calculations, we obtain the results.  $\square$

#### 4. Non-Hermitian $\mathbb{R}$ -Complex Finsler Space with an Arctangent Metric

In this section, we deal with the non-Hermitian  $\mathbb{R}$ -Complex Finsler space with an arctangent metric given in equation (1.1).

For the non-Hermitian  $\mathbb{R}$ -Complex Finsler space ( $a_{i\bar{j}} = 0$ ), we use the following abbreviations:

$$\begin{aligned}
 l_i &= a_{ij} \eta^j, \gamma = a_{jk} \eta^j \eta^k = l_k \eta^k, \theta = b_j \eta^j, \omega = b_j b^j, \\
 b^k &= a^{jk} b_j, b_l = b^k a_{kl}, \delta = a_{jk} \eta^j b^k = l_k b^k, l^j = a^{ji} l_i = \eta^j. \quad (4.1)
 \end{aligned}$$

**Theorem 4.1.** For a non-Hermitian  $\mathbb{R}$ -Complex Finsler space  $(M, F)$ , where  $F$  is an arctangent metric, we have

i) the contravariant tensor  $g^{ji}$  which is given in equation

$$g^{ji} = \frac{1}{\rho_0} \left[ a^{ji} + \left\{ \frac{t_1}{\tau_1} - \frac{\theta^2 t_1^2 t_2}{\tau_1^2 \tau^2} \right\} \eta^i \eta^j - \frac{t_2 b^i b^j}{\tau_2} - \frac{\theta t_1 t_2 (b^i \eta^j + b^j \eta^i)}{\tau_1 \tau_2} - \frac{A^2 \eta^i \eta^j + AB(b^i \eta^j + b^j \eta^i) + B^2 b^i b^j}{\tau_3} \right]. \quad (4.2)$$

ii) The  $\det(g_{ij})$  which is given in equation

$$\det(g_{ij}) = (\rho_0)^n \tau_1 \tau_2 \tau_3 \det(a_{ij}), \quad (4.3)$$

where  $\rho_0, t_1$  and  $t_2$  are given in equations (3.5) and (3.10), rest terms are

$$\begin{aligned} A &= \left\{ 1 + \frac{t_1}{\tau_1} - \frac{\theta^2 t_1^2 t_2}{(\tau_1)^2 \tau_2} \right\} \gamma - \frac{\theta t_1 t_2}{(\tau_1)^3 \tau_2}, \\ B &= -\frac{t_2 \theta}{\tau_2} - \frac{\theta t_1 t_2 \gamma}{\tau_1 \tau_2}, \\ \tau_1 &= 1 - t_1 \gamma, \\ \tau_2 &= 1 + t_2 \left( \omega + \frac{t_1 \theta^2}{\tau_1} \right), \\ \tau_3 &= 1 + (A\gamma + B\theta) \sqrt{t_3}. \end{aligned} \quad (4.4)$$

*Proof.* Now, apply proposition 2.4 to  $g_{ij}$  in equation (3.8) and follow the steps:

**Step 1.** [Suppose  $Q_{ij} = a_{ij}$  and  $C_i = \sqrt{t_1} l_i$ ]

From our assumption, we get

$$Q^{ji} = a^{ji}$$

and

$$C^2 = C_i C^i = \sqrt{t_1} l_i \times Q^{ji} \times C_j = \sqrt{t_1} l_i \times a^{ji} \times \sqrt{t_1} l_j = t_1 \times l_i a^{ji} l_j = t_1 \gamma.$$

By applying proposition 2.4, we get

$$\det(H_{ij}) = \det(a_{ij} - t_1 l_i l_j) = (1 - t_1 \gamma) \det(a_{ij}) = \tau_1 \det(a_{ij}), \quad (4.5)$$

and, for  $\tau_1 = 1 - t_1 \gamma \neq 0$ ,  $(H_{ij}) = (a_{ij} - t_1 l_i l_j)$  is invertible and its inverse is given by:

$$H^{ji} = a^{ji} + \frac{t_1 \eta^i \eta^j}{\tau_1}. \quad (4.6)$$

**Step 2.** [Suppose  $Q_{ij} = a_{ij} - t_1 l_i l_j$  and  $C_i = \sqrt{t_2} b_i$ ]

Using the equations (4.1), (4.6), and our supposition, we get

$$Q^{ji} = a^{ji} + \frac{t_1 \eta^i \eta^j}{\tau_1}.$$



Using the previous equation, we get

$$\begin{aligned} C^i &= Q^{ji}C_j = \left( a^{ji} + \frac{t_1\eta^i\eta^j}{\tau_1} \right) \sqrt{t_2}b_j \\ &= \left( b^i + \frac{t_1\theta\eta^i}{\tau_1} \right) \sqrt{t_2}, \end{aligned}$$

which implies

$$C^2 = t_2 \left( \omega + \frac{t_1\theta^2}{\tau_1} \right),$$

and

$$1 + C^2 = 1 + t_2 \left( \omega + \frac{t_1\theta^2}{\tau_1} \right) = \tau_2(\text{say}).$$

Now, by applying proposition 2.4, we get

$$\det(H_{ij}) = \det(a_{ij} - t_1l_il_j + t_2b_ib_j) = \tau_1\tau_2\det(a_{ij}), \quad (4.7)$$

and, for  $\tau_2$  and  $\tau_1 \neq 0$ , the inverse of  $(H_{ij}) = (a_{ij} - t_1l_il_j + t_2b_ib_j)$  exists and it is

$$\begin{aligned} H^{ji} &= a^{ji} + \left\{ \frac{t_1}{\tau_1} - \frac{\theta^2 t_1^2 t_2}{(\tau_1)^2 \tau_2} \right\} \eta^i \eta^j - \frac{t_2 b^i b^j}{\tau_2} \\ &\quad + \frac{\theta t_1 t_2 (b^i \eta^j + b^j \eta^i)}{\tau_1 \tau_2}. \end{aligned} \quad (4.8)$$

**Step 3.** [Suppose  $Q_{ij} = a_{ij} - t_1l_il_j + t_2b_ib_j$  and  $C_i = \sqrt{t_3}\eta_i$ ] Using the equation (4.8) and our supposition, we get

$$\begin{aligned} Q^{ji} &= a^{ji} + \left\{ \frac{t_1}{\tau_1} - \frac{\theta^2 t_1^2 t_2}{(\tau_1)^2 \tau_2} \right\} \eta^i \eta^j - \frac{t_2 b^i b^j}{\tau_2} \\ &\quad + \frac{\theta t_1 t_2 (b^i \eta^j + b^j \eta^i)}{\tau_1 \tau_2}. \end{aligned}$$

Using the previous equation, we get

$$C^i = A\eta^i + Bb^i,$$

where

$$\begin{aligned} A &= \left\{ 1 + \frac{t_1}{\tau_1} - \frac{\theta^2 t_1^2 t_2}{(\tau_1)^2 \tau_2} \right\} \gamma - \frac{\theta t_1 t_2}{(\tau_1)^3 \tau_2}, \\ B &= -\frac{t_2 \theta}{\tau_2} - \frac{\theta t_1 t_2 \gamma}{\tau_1 \tau_2}, \end{aligned} \quad (4.9)$$

which implies

$$C^2 = Q^{ji}C_j = (A\gamma + B\theta)\sqrt{t_3}, \quad 1 + C^2 = 1 + (A\gamma + B\theta)\sqrt{t_3} = \tau_3(\text{say}),$$

and

$$C^i C^j = A^2 \eta^i \eta^j + AB(b^i \eta^j + b^j \eta^i) + B^2 b^i b^j.$$

Now, by using proposition 2.4, we get

$$\det(H_{ij}) = \det(a_{ij} - t_1 l_i l_j + t_2 b_i b_j + t_3 \eta_i \eta_j) = \tau_1 \tau_2 \tau_3 \det(a_{ij}), \quad (4.10)$$

and for non-zero  $\tau_i (i = 1, 2, 3, )$ , the inverse of  $(H_{ij}) = (a_{ij} - t_1 l_i l_j + t_2 b_i b_j + t_3 \eta_i \eta_j)$  exists and it is

$$H^{ji} = a^{ji} + \left\{ \frac{t_1}{\tau_1} - \frac{\theta^2 t_1^2 t_2}{\tau_1^2 \tau^2} \right\} \eta^i \eta^j - \frac{t_2 b^i b^j}{\tau_2} - \frac{\theta t_1 t_2 (b^i \eta^j + b^j \eta^i)}{\tau_1 \tau_2} - \frac{A^2 \eta^i \eta^j + AB(b^i \eta^j + b^j \eta^i) + B^2 b^i b^j}{\tau_3}. \quad (4.11)$$

But  $g_{ij} = \rho_0 H_{ij}$ , where  $H_{ij}$  is given in the previous equation. Thus,

$$g^{ji} = \frac{1}{\rho_0} H^{ji}$$

and

$$\det(g_{ij}) = (\rho_0)^n \det(H_{ij}).$$

Using the equations (4.10) and (4.11), we get

$$g^{ji} = \frac{1}{\rho_0} \left[ a^{ji} + \left\{ \frac{t_1}{\tau_1} - \frac{\theta^2 t_1^2 t_2}{\tau_1^2 \tau^2} \right\} \eta^i \eta^j - \frac{t_2 b^i b^j}{\tau_2} - \frac{\theta t_1 t_2 (b^i \eta^j + b^j \eta^i)}{\tau_1 \tau_2} - \frac{A^2 \eta^i \eta^j + AB(b^i \eta^j + b^j \eta^i) + B^2 b^i b^j}{\tau_3} \right], \quad (4.12)$$

and

$$\det(g_{ij}) = (\rho_0)^n \tau_1 \tau_2 \tau_3 \det(a_{ij}). \quad (4.13)$$

Hence the statement holds.  $\square$

Now, in a non-Hermitian  $\mathbb{R}$ -Complex Finsler space  $(M, F)$ , where  $F$  is an arctangent metric, we have the following properties:

$$\gamma + \bar{\gamma} = l_i \eta^i + l_{\bar{j}} \bar{\eta}^{\bar{j}} = a_{ij} \eta^j \eta^i + a_{\bar{j}\bar{k}} \bar{\eta}^{\bar{k}} \bar{\eta}^{\bar{j}} = 2\alpha^2, \quad (4.14)$$

$$\theta + \bar{\theta} = b_j \eta^j + b_{\bar{j}} \bar{\eta}^{\bar{j}} = 2\beta, \quad \delta = \theta. \quad (4.15)$$

**Proposition 4.2.** *Let us consider a non-Hermitian  $\mathbb{R}$ -Complex Finsler space  $(M, F)$ , where  $F$  is an arctangent metric. This space satisfies the properties given in equations (4.14) and (4.15).*

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